

# Hartreeho-Fockova metoda

(aneb začíná přituhovat)

**Kvantová chemie**

**Lekce 9**

# Osnova

1. Ingredience Hartreeho-Fockovy metody
2. Hartreeho-Fockova metoda
3. Hartreeho-Fockovy rovnice
4. Energie
5. Restringovaná Hartreeho-Fockova metoda

# Ingredience H-F metody

## Ritzův variační princip

$$E_0 = \min_{\|\psi\|=1} \langle \psi | \hat{H} | \psi \rangle \quad (\text{počítáme tedy „jen“ energii základního stavu})$$

## Hamiltonův operátor

(X-representace, B-O aproximace, elektrostatické přiblížení)

$$\hat{H} \equiv \hat{H}_e = \sum_{k=1}^n \left( -\frac{\hbar^2}{2m_e} \Delta_k \right) + \sum_{J=1}^{N-1} \sum_{K=J+1}^N \frac{Z_J Z_K \tilde{e}^2}{\|\vec{R}_J - \vec{R}_K\|} + \sum_{j=1}^{n-1} \sum_{k=j+1}^n \frac{\tilde{e}^2}{\|\vec{r}_j - \vec{r}_k\|} - \sum_{J=1}^N \sum_{j=1}^n \frac{Z_J \tilde{e}^2}{\|\vec{R}_J - \vec{r}_j\|}$$

## Vlnová funkce

(Hartreeho-Fockova aproximace)

$$\psi = \frac{1}{\sqrt{n!}} \det \begin{pmatrix} \varphi_1(1) & \cdots & \varphi_n(1) \\ \vdots & \ddots & \vdots \\ \varphi_1(n) & \cdots & \varphi_n(n) \end{pmatrix}, \quad \langle \varphi_j | \varphi_k \rangle \equiv \sum_{\xi=-1/2}^{+1/2} \int_{\mathbb{R}^3} \varphi_j^*(\vec{r}, \xi) \varphi_k(\vec{r}, \xi) d\vec{r} = \delta_{jk}$$

# Ingredience H-F metody

## Ritzův variační princip

$$E_0 = \min_{\|\psi\|=1} \langle \psi | \hat{H} | \psi \rangle \quad (\text{počítáme tedy „jen“ energii základního stavu})$$

## Hamiltonův operátor

(X-representace, B-O aproximace, elektrostatické přiblížení)

$$\tilde{e}^2 = \frac{e^2}{4\pi\epsilon_0}$$

$$\hat{H} = \hat{H}_e = \sum_{k=1}^n \left( -\frac{\hbar^2}{2m_e} \Delta_k \right) + \sum_{J=1}^{N-1} \sum_{K=J+1}^N \frac{Z_J Z_K \tilde{e}^2}{\|\vec{R}_J - \vec{R}_K\|} + \sum_{j=1}^{n-1} \sum_{k=j+1}^n \frac{\tilde{e}^2}{\|\vec{r}_j - \vec{r}_k\|} - \sum_{J=1}^N \sum_{j=1}^n \frac{Z_J \tilde{e}^2}{\|\vec{R}_J - \vec{r}_j\|}$$

## Vlnová funkce

(Hartreeho-Fockova aproximace)

$$1 = \vec{r}_1, \xi_1 \text{ atd.}$$

$$\psi = \frac{1}{\sqrt{n!}} \det \begin{pmatrix} \varphi_1(1) & \cdots & \varphi_n(1) \\ \vdots & \ddots & \vdots \\ \varphi_1(n) & \cdots & \varphi_n(n) \end{pmatrix}, \quad \langle \varphi_j | \varphi_k \rangle \equiv \sum_{\xi=-1/2}^{+1/2} \int_{\mathbb{R}^3} \varphi_j^*(\vec{r}, \xi) \varphi_k(\vec{r}, \xi) d\vec{r} = \delta_{jk}$$

# H-F metoda

## Suma sumárum

$$E_0 = \min_{\varphi_1, \dots, \varphi_n} \langle \psi(\varphi_1, \dots, \varphi_n) | \hat{H} | \psi(\varphi_1, \dots, \varphi_n) \rangle, \text{ s vazbami } \langle \varphi_j | \varphi_k \rangle = \delta_{jk}$$

## Metoda Lagrangeových multiplikátorů

$$E_0 = \min_{\varphi_1, \dots, \varphi_n; \varepsilon_{jk}} \{ \langle \psi(\varphi_1, \dots, \varphi_n) | \hat{H} | \psi(\varphi_1, \dots, \varphi_n) \rangle - \sum_{j,k} \varepsilon_{jk} (\langle \varphi_j | \varphi_k \rangle - \delta_{jk}) \}$$

# H-F metoda

## Suma sumárum

$$E_0 = \min_{\varphi_1, \dots, \varphi_n} \langle \psi(\varphi_1, \dots, \varphi_n) | \hat{H} | \psi(\varphi_1, \dots, \varphi_n) \rangle, \text{ s vazbami } \langle \varphi_i | \varphi_k \rangle = \delta_{ik}$$

$\varepsilon = \{\varepsilon_{jk}\}$  je hermitovská (symetrická) matice

$$\mathbf{U}^+ \varepsilon \mathbf{U} = \text{diag}\{\tilde{\varepsilon}_1, \dots, \tilde{\varepsilon}_n\} \quad (\mathbf{Q}^T \varepsilon \mathbf{Q} = \text{diag}\{\tilde{\varepsilon}_1, \dots, \tilde{\varepsilon}_n\})$$

## Metoda Lagrangeových multiplikátorů

$$E_0 = \min_{\varphi_1, \dots, \varphi_n; \varepsilon_{jk}} \left\{ \langle \psi(\varphi_1, \dots, \varphi_n) | \hat{H} | \psi(\varphi_1, \dots, \varphi_n) \rangle - \sum_{j,k} \varepsilon_{jk} (\langle \varphi_j | \varphi_k \rangle - \delta_{jk}) \right\}$$

# H-F metoda

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$$E_0 = \min_{\varphi_1, \dots, \varphi_n} \langle \psi(\varphi_1, \dots, \varphi_n) | \hat{H} | \psi(\varphi_1, \dots, \varphi_n) \rangle, \text{ s vazbami } \langle \varphi_i | \varphi_l \rangle = \delta_{il}$$

$\varepsilon = \{\varepsilon_{jk}\}$  je hermitovská (symetrická) matice  
 $\mathbf{U}^+ \varepsilon \mathbf{U} = \text{diag}\{\tilde{\varepsilon}_1, \dots, \tilde{\varepsilon}_n\}$  ( $\mathbf{Q}^T \varepsilon \mathbf{Q} = \text{diag}\{\tilde{\varepsilon}_1, \dots, \tilde{\varepsilon}_n\}$ )

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$|\varphi_k\rangle \rightarrow |\tilde{\varphi}_k\rangle = \sum_{l=1}^n U_{lk}^* |\varphi_l\rangle$  ( $|\varphi_k\rangle \rightarrow |\tilde{\varphi}_k\rangle = \sum_{l=1}^n Q_{lk} |\varphi_l\rangle$ ), a vynecháme vlnovky

# H-F metoda

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Me

$$\begin{aligned} \sum_{j,k} \varepsilon_{jk} \langle \varphi_j | \varphi_k \rangle &= \boldsymbol{\varphi}^+ \boldsymbol{\varepsilon} \boldsymbol{\varphi} = \boldsymbol{\varphi}^+ \mathbf{U} \mathbf{U}^+ \boldsymbol{\varepsilon} \mathbf{U} \mathbf{U}^+ \boldsymbol{\varphi} = \\ &= (\boldsymbol{\varphi}^+ \mathbf{U}) (\mathbf{U}^+ \boldsymbol{\varepsilon} \mathbf{U}) (\mathbf{U}^+ \boldsymbol{\varphi}) = (\mathbf{U}^+ \boldsymbol{\varphi})^+ (\mathbf{U}^+ \boldsymbol{\varepsilon} \mathbf{U}) (\mathbf{U}^+ \boldsymbol{\varphi}) \end{aligned}$$

ú

$$\langle \psi(\varphi_1, \dots, \varphi_n) | \psi(\varphi_1, \dots, \varphi_n) \rangle - \sum_{j,k} \varepsilon_{jk} (\langle \varphi_j | \varphi_k \rangle - \delta_{jk})$$

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# H-F metoda

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$$\langle \psi(\varphi_1, \dots, \varphi_n) | \psi(\varphi_1, \dots, \varphi_n) \rangle - \sum_{j,k} \varepsilon_{jk} (\langle \varphi_j | \varphi_k \rangle - \delta_{jk})$$

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$$E_0 = \min_{\varphi_k; \varepsilon_k} \{ \langle \psi(\varphi_1, \dots, \varphi_n) | \hat{H} | \psi(\varphi_1, \dots, \varphi_n) \rangle - \sum_k \varepsilon_k \langle \varphi_k | \varphi_k \rangle \}$$

# H-F metoda

## Funkcionál energie

$$\langle \psi | \hat{H} | \psi \rangle = \langle \psi | \sum_{k=1}^n \left( -\frac{\hbar^2}{2m_e} \Delta_k - \sum_{J=1}^N \frac{Z_J \tilde{e}^2}{\|\vec{R}_J - \vec{r}_k\|} \right) | \psi \rangle + \quad (\text{jednoelektronový příspěvek})$$

$$\langle \psi | \sum_{J=1}^{N-1} \sum_{K=J+1}^N \frac{Z_J Z_K \tilde{e}^2}{\|\vec{R}_J - \vec{R}_K\|} | \psi \rangle + \quad (\text{příspěvek elstat repulze jader})$$

$$\frac{1}{2} \langle \psi | \sum_{j,k=1, j \neq k}^n \frac{\tilde{e}^2}{\|\vec{r}_j - \vec{r}_k\|} | \psi \rangle = \quad (\text{dvojelektronový příspěvek})$$

$$\sum_{k=1}^n \langle \varphi_k | \left( -\frac{\hbar^2}{2m_e} \Delta_k - \sum_{J=1}^N \frac{Z_J \tilde{e}^2}{\|\vec{R}_J - \vec{r}_k\|} \right) | \varphi_k \rangle +$$

$$\sum_{J=1}^{N-1} \sum_{K=J+1}^N \frac{Z_J Z_K \tilde{e}^2}{\|\vec{R}_J - \vec{R}_K\|} +$$

$$\frac{1}{2} \sum_{j,k=1, j \neq k}^n \langle \varphi_j \varphi_k | \frac{\tilde{e}^2}{\|\vec{r}_j - \vec{r}_k\|} | \varphi_j \varphi_k \rangle - \frac{1}{2} \sum_{j=1, j \neq k}^n \langle \varphi_j \varphi_k | \frac{\tilde{e}^2}{\|\vec{r}_j - \vec{r}_k\|} | \varphi_k \varphi_j \rangle$$

# H-F metoda

## Funkcionál energie

$$\langle \psi | \hat{H} | \psi \rangle = \langle \psi | \sum_{k=1}^n \left( -\frac{\hbar^2}{2m_e} \Delta_k - \sum_{J=1}^{N-1} \frac{Z_J \tilde{e}^2}{\|\vec{R}_J - \vec{r}_k\|} \right) | \psi \rangle + \langle \psi | \sum_{J=1}^{N-1} \sum_{K=J+1}^N \frac{Z_J Z_K \tilde{e}^2}{\|\vec{R}_J - \vec{R}_K\|} | \psi \rangle + \frac{1}{2} \langle \psi | \sum_{j,k=1, j \neq k}^n \frac{\tilde{e}^2}{\|\vec{r}_j - \vec{r}_k\|} | \psi \rangle =$$

(jednoelektronový příspěvek)

(příspěvek elstat repulze jader)

(dvojelektronový příspěvek)

$$\sum_{k=1}^n \left\langle \varphi_k \left| \left( -\frac{\hbar^2}{2m_e} \Delta_k - \sum_{J=1}^N \frac{Z_J \tilde{e}^2}{\|\vec{R}_J - \vec{r}_k\|} \right) \right| \varphi_k \right\rangle +$$

$$\sum_{J=1}^{N-1} \sum_{K=J+1}^N \frac{Z_J Z_K \tilde{e}^2}{\|\vec{R}_J - \vec{R}_K\|} +$$

$$\frac{1}{2} \sum_{j,k=1, j \neq k}^n \left\langle \varphi_j \varphi_k \left| \frac{\tilde{e}^2}{\|\vec{r}_j - \vec{r}_k\|} \right| \varphi_j \varphi_k \right\rangle - \frac{1}{2} \sum_{j=1, j \neq k}^n \left\langle \varphi_j \varphi_k \left| \frac{\tilde{e}^2}{\|\vec{r}_j - \vec{r}_k\|} \right| \varphi_k \varphi_j \right\rangle$$

# H-F rovnice

## Matematická podstata H-F metody

- úloha z oblasti variačního počtu (s omezeními)
- Lagrangeovy rovnice  $\rightarrow$  ...

## Hartreeho-Fockovy rovnice

$$\left\{ -\frac{\hbar^2}{2m_e} \Delta - \sum_{J=1}^N \frac{Z_J \tilde{e}^2}{\|\vec{r} - \vec{R}_J\|} \right\} \varphi_k(\vec{r}, \xi) + \left\{ \sum_{j=1, j \neq k}^n \int_{\mathbb{R}^3} \frac{\tilde{e}^2}{\|\vec{r} - \vec{r}'\|} [\sum_{\xi'} \varphi_j^*(\vec{r}', \xi') \varphi_j(\vec{r}', \xi')] d\vec{r}' \right\} \varphi_k(\vec{r}, \xi) -$$
$$- \sum_{j=1, j \neq k}^n \left\{ \int_{\mathbb{R}^3} \frac{\tilde{e}^2}{\|\vec{r} - \vec{r}'\|} [\sum_{\xi'} \varphi_j^*(\vec{r}', \xi') \varphi_k(\vec{r}', \xi')] d\vec{r}' \right\} \varphi_j(\vec{r}, \xi) = \varepsilon_k \varphi_k(\vec{r}, \xi)$$

$$k = 1, \dots, n$$

# H-F rovnice

## Interpretace

- matematická
  - soustava  $n$  vázaných integro-diferenciálních rovnic
  - pro  $n$  (reálných) funkcí 4 (3+1) proměnných  $\varphi_k(\vec{r}, \xi)$
- fyzikální
  - **jednoelektronové Schrödingerovy rovnice**

$$\left\{ -\frac{\hbar^2}{2m_e} \Delta - \sum_{J=1}^N \frac{Z_J \tilde{e}^2}{\|\vec{r} - \vec{R}_J\|} \right\} \varphi_k(\vec{r}, \xi) + \left\{ \sum_{j=1, j \neq k}^n \int_{\mathbb{R}^3} \frac{\tilde{e}^2}{\|\vec{r} - \vec{r}'\|} [\sum_{\xi'} \varphi_j^*(\vec{r}', \xi') \varphi_j(\vec{r}', \xi')] d\vec{r}' \right\} \varphi_k(\vec{r}, \xi) -$$
$$- \sum_{j=1, j \neq k}^n \left\{ \int_{\mathbb{R}^3} \frac{\tilde{e}^2}{\|\vec{r} - \vec{r}'\|} [\sum_{\xi'} \varphi_j^*(\vec{r}', \xi') \varphi_k(\vec{r}', \xi')] d\vec{r}' \right\} \varphi_j(\vec{r}, \xi) = \varepsilon_k \varphi_k(\vec{r}, \xi)$$

# H-F rovnice

## Interpretace

- matematická
  - soustava  $n$  vázaných integro-diferenciálních rovnic
  - pro  $n$  (reálných) funkcí 4 (3+1) proměnných  $\varphi_k(\vec{r}, \xi)$
- fyzikální

operátor kinetické energie elektronu  $k$

interakce elektronu  $k$  se všemi jádry

interakce elektronu  $k$  s ostatními elektrony

$$\left\{ -\frac{\hbar^2}{2m_e} \Delta - \sum_{J=1}^N \frac{Z_J \tilde{e}^2}{\|\vec{r} - \vec{R}_J\|} \right\} \varphi_k(\vec{r}, \xi) + \left\{ \sum_{j=1, j \neq k}^n \int_{\mathbb{R}^3} \frac{\tilde{e}^2}{\|\vec{r} - \vec{r}'\|} \left[ \sum_{\xi'} \varphi_j^*(\vec{r}', \xi') \varphi_j(\vec{r}', \xi') \right] d\vec{r}' \right\} \varphi_k(\vec{r}, \xi) - \sum_{j=1, j \neq k}^n \left\{ \int_{\mathbb{R}^3} \frac{\tilde{e}^2}{\|\vec{r} - \vec{r}'\|} \left[ \sum_{\xi'} \varphi_j^*(\vec{r}', \xi') \varphi_k(\vec{r}', \xi') \right] d\vec{r}' \right\} \varphi_j(\vec{r}, \xi) = \varepsilon_k \varphi_k(\vec{r}, \xi)$$

výměnná interakce

# H-F rovnice

## Hartreeho metoda (rovnice)

- vlnová funkce

$$\psi(1,2, \dots, n) = \varphi_1(1)\varphi_2(2) \dots \varphi_n(n)$$

- Hartreeho rovnice

$$\left\{ -\frac{\hbar^2}{2m_e} \Delta - \sum_{J=1}^N \frac{Z_J \tilde{e}^2}{\|\vec{r} - \vec{R}_J\|} \right\} \varphi_k(\vec{r}, \xi) + \left\{ \sum_{j=1, j \neq k}^n \int_{\mathbb{R}^3} \frac{\tilde{e}^2}{\|\vec{r} - \vec{r}'\|} [\sum_{\xi'} \varphi_j^*(\vec{r}', \xi') \varphi_j(\vec{r}', \xi')] d\vec{r}' \right\} \varphi_k(\vec{r}, \xi) = \varepsilon_k \varphi_k(\vec{r}, \xi)$$

chybí výměnný člen:  $-\sum_{j=1, j \neq k}^n \left\{ \int_{\mathbb{R}^3} \frac{\tilde{e}^2}{\|\vec{r} - \vec{r}'\|} [\sum_{\xi'} \varphi_j^*(\vec{r}', \xi') \varphi_k(\vec{r}', \xi')] d\vec{r}' \right\} \varphi_j(\vec{r}, \xi)$

# Energie (základního stavu)

## Variační princip

$$E_0 = \min_{\|\psi\|=1} \langle \psi | \hat{H} | \psi \rangle \leq \langle \psi_{\text{HF}} | \hat{H} | \psi_{\text{HF}} \rangle$$

## Hartreeho-Fockovo přiblížení

$$E_0 \approx E_{\text{HF}} \equiv \langle \psi_{\text{HF}} | \hat{H} | \psi_{\text{HF}} \rangle \neq \sum_{k=1}^n \varepsilon_k$$

## Proč?

- chybí odpuzivá interakce jader:  $\sum_{J=1}^{N-1} \sum_{K=J+1}^N \frac{Z_J Z_K \tilde{e}^2}{\|\vec{R}_J - \vec{R}_K\|}$
- vzájemná interakce elektronů je započtena dvakrát

$$E_{\text{HF}} = \sum_{k=1}^n \varepsilon_k + \sum_{J=1}^{N-1} \sum_{K=J+1}^N \frac{Z_J Z_K \tilde{e}^2}{\|\vec{R}_J - \vec{R}_K\|} - \left\langle \psi_{\text{HF}} \left| \sum_{j=1}^{n-1} \sum_{k=j+1}^n \frac{\tilde{e}^2}{\|\vec{r}_j - \vec{r}_k\|} \right| \psi_{\text{HF}} \right\rangle$$



# Restringovaná H-F metoda (uzavřené slupky)

## Vlnová funkce

- sudý počet elektronů, po páru s opačnou orientací projekce spinu
- jednoelektronové funkce
  - $\varphi_1(\vec{r}, \xi) = \phi_1(\vec{r})\alpha(\xi)$   
 $\varphi_2(\vec{r}, \xi) = \phi_1(\vec{r})\beta(\xi)$
  - $\varphi_3(\vec{r}, \xi) = \phi_2(\vec{r})\alpha(\xi)$   
 $\varphi_4(\vec{r}, \xi) = \phi_2(\vec{r})\beta(\xi)$
  - ...
  - $\varphi_{n-1}(\vec{r}, \xi) = \phi_{n/2}(\vec{r})\alpha(\xi)$   
 $\varphi_n(\vec{r}, \xi) = \phi_{n/2}(\vec{r})\beta(\xi)$

# Restringovaná H-F metoda (uzavřené slupky)

## Vlnová funkce

$$\alpha\left(+\frac{1}{2}\right) = 1, \alpha\left(-\frac{1}{2}\right) = 0$$

$$\beta\left(+\frac{1}{2}\right) = 0, \beta\left(-\frac{1}{2}\right) = 1$$

- sudý počet elektronů  $n$  – reprezentací projekce spinu

- jednoelektronové orbitály
  - $\varphi_1(\vec{r}, \xi) = \phi_1(\vec{r})\alpha(\xi)$   
 $\varphi_2(\vec{r}, \xi) = \phi_1(\vec{r})\beta(\xi)$
  - $\varphi_3(\vec{r}, \xi) = \phi_2(\vec{r})\alpha(\xi)$   
 $\varphi_4(\vec{r}, \xi) = \phi_2(\vec{r})\beta(\xi)$
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## Hartreeho-Fockovy rovnice

$$\left\{ -\frac{\hbar^2}{2m_e} \Delta - \sum_{J=1}^N \frac{Z_J \tilde{e}^2}{\|\vec{r} - \vec{R}_J\|} \right\} \phi_k(\vec{r}) + 2 \left\{ \sum_{j=1, j \neq k}^{n/2} \int_{\mathbb{R}^3} \frac{\tilde{e}^2}{\|\vec{r} - \vec{r}'\|} \phi_j^*(\vec{r}') \phi_j(\vec{r}') d\vec{r}' \right\} \phi_k(\vec{r}) -$$
$$- \sum_{j=1, j \neq k}^{n/2} \left\{ \int_{\mathbb{R}^3} \frac{\tilde{e}^2}{\|\vec{r} - \vec{r}'\|} \phi_j^*(\vec{r}') \phi_k(\vec{r}') d\vec{r}' \right\} \phi_j(\vec{r}) = \varepsilon_k \phi_k(\vec{r})$$

$$k = 1, \dots, n/2$$

**Konec lekce 9.**